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Yuri P. Panarin ^a, Fouad M. Aliev ^a & Charles Rosenblatt ^b

^a Department of Physics, University of Puerto Rico, P.O. Box 23343, PR, USA

^b Department of Physics, Case Western Reserve University, Cleveland, OH, USA

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The Appearance of Ferrielectric Phases in Confined Liquid Crystal Studied by Photon Correlation Spectroscopy

YURI P. PANARIN^a, FOUAD M. ALIEV^a and
CHARLES ROSENBLATT^b

^a*Department of Physics, University of Puerto Rico, P.O. Box 23343, PR, USA; and*

^b*Department of Physics, Case Western Reserve University, Cleveland, OH, USA*

The influence of confinement on the dynamic properties and appearance of ferrielectric phases has been investigated by dynamic light scattering, performed in the bulk AFLC as well as confined in 200 Å and 2000 Å cylindrical pores. The analysis of the experimental results shows that the ferrielectric phases are not formed in 200 Å pores, while they appear in 2000 Å pores as well as in bulk.

Keywords: antiferroelectric liquid crystal; photon correlation spectroscopy

INTRODUCTION

The appearance of ferrielectric phases in antiferroelectric liquid crystals (AFLC) can be understood as a result of the competition between the antiferro- and ferroelectric interactions in adjacent smectic layers that stabilize the SmC_A and the SmC^* phases. In ferrielectric liquid crystals the directions of the local spontaneous polarizations in any two neighboring smectic layers can be ferroelectric-like or antiferroelectric-like. The antiferroelectric order is believed to exist due to steric factors since it appears in racemate chiral liquid

crystals^[1] as well as in non-chiral tilted smectics^[2]. On the other hand we are unaware of any report of ferrielectricity in non-chiral and/or racemic liquid crystals. It seems, therefore, that the existence of local spontaneous polarization is a determining factor for ferrielectricity. The Ising model^[3], which takes into account the long-range electrostatic interaction between smectic layers, predicts the existence of an infinite number of ferrielectric phases^[1]. According to the Ising model, these phases can be characterized by the parameter $q_T = F/(A+F)$, which denotes the fraction of ferroelectric ordering (F) which appears together with antiferroelectric ordering (A) in the periodic structure. Following this definition, the SmC_A is the antiferroelectric phase with $q_T=0$, the SmC_γ is ferrielectric phase with $q_T=1/3$, the AF is antiferroelectric phase with $q_T=1/2$, and the SmC^* is ferroelectric with $q_T=1$. Additionally, other ferrielectric subphases (FiLC , SmC_α) have been found, which cannot be described by the q_T parameter.

The Ising model was developed for bulk liquid crystals; therefore the question "What happens to ferrielectric phases in confined liquid crystals?" is of great interest and importance for understanding the nature of ferrielectricity in liquid crystals. In parallel cylindrical pores the molecular tilt directions are equivalent with respect to azimuthal angle (if planar boundary conditions are assumed). Therefore in cylindrical pores the effect of the surface interaction can be ignored and we can investigate the effect of confinement on the appearance of ferrielectric phases.

Dynamic properties of bulk antiferroelectric and ferroelectric phases have been investigated by dielectric^[4-10], electrooptic^[5,10] and photon correlation^[11-16] spectroscopy. Summarizing the results of these investigations, two collective relaxation processes have been observed in antiferroelectric phase:

(i) antiphase mode^[7] (or optic-like phason^[15]) and (ii) in-phase mode^[7] (or sound-like phason^[14]).

Investigation of AFLC confined in cylindrical pores is rather difficult, because of specific confining geometry. In dielectric spectroscopy experiment, applied to our system the direction of electric field is along the smectic layer normal and the information about collective modes is missed. Polarizing microscopy is not effective because the pore diameter is smaller than light wavelength. Therefore, in this paper we report on the influence of confinement on the appearance of antiferroelectric/ ferroelectric phases in AFLC investigated by photon correlation spectroscopy.

Dynamic light scattering is a powerful technique to investigate dynamic properties of liquid crystals both in bulk and confined geometry. In dynamic light scattering the intensity/intensity autocorrelation function provides information about relaxation processes in the sample, which are different in different ferroelectric/ antiferroelectric phases.

EXPERIMENTAL

We have performed investigations in bulk AFLC and confined in cylindrical pores of Anopore membranes with different diameters: 200 Å and 2000 Å. The AFLC used in the experiments was (R) - 1 - methylheptyl - 4 -(4''- n-dodecyloxybiphenyl-4'-carbonoyloxy) - 3 - fluorobenzoate (AS-573, Hull, UK).

This AFLC is of great interest because it possesses a wide variety of different antiferroelectric and ferroelectric. The phase transition sequence of the bulk sample has been determined by the macroscopic polarization measurements and dielectric spectroscopy^[9]:



In our experiments the bulk liquid crystal was placed between two untreated glass plates having a 100 μm gap. The texture of liquid crystal was observed to be planar disordered with a helix. The helical pitch was about 0.2–0.3 μm (less than a wavelength of visible light λ). For the investigations in the confined geometry we used Anopore membranes of thickness 60 μm . The liquid crystal structure within the pores is believed to be helical, with helical axes along the pores, because a selective reflection in blue region was observed.

We performed photon correlation spectroscopy using a He-Ne laser operating at wavelength 6328Å, and an ALV-5000/Fast Digital Multiple Tau Correlator (real time) operating over delay times from 12.5 ns up to 10^3 s. The depolarized component of scattered light was investigated. All the data discussed below were obtained at $\theta=30^\circ$. The data were taken after stabilizing the temperature for one hour at each value of T , sufficiently long to guarantee thermal equilibrium.

RESULTS

The temperature-dependence measurements were carried out in the temperature range $72^\circ\text{C} < T < 94^\circ\text{C}$. Figure 1(a-c) presents the intensity/intensity autocorrelation functions of AS573 in the 100 μm cell (bulk) and in pores for three different temperatures corresponding to different phases over the time range $10^{-2} < \tau < 10^4$ ms.

The autocorrelation functions of the bulk AFLC consist of (at most) three decays corresponding to three relaxation processes (Fig. 1b, curve 1): (i) the “fast” process ($\tau = 27 \mu\text{s}$), (ii) the “middle frequency” process ($\tau = 50$ ms) and (iii) the “low frequency” process ($\tau = 850$ ms). From these figures we observe

some similarities and differences in the shape of autocorrelation functions obtained in bulk and in confined liquid crystal.

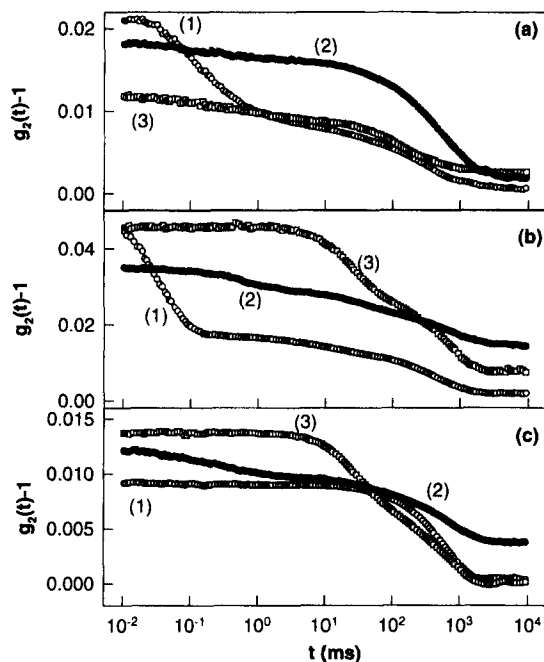


FIGURE 1 Intensity/intensity autocorrelation functions of AS573 in bulk (1), in 2000 Å pores (2), and in 200 Å pores (3) measured at (a) 75.5°C, (b) 85.5°C, and (c) 91.5°C.

The amplitude of the fast relaxation process strongly depends on the temperature. In the bulk sample this process is very prominent in the temperature range below 88°C, which corresponds to antiferroelectric/ferrielectric phases (Fig.1a,b, curve 1). In the temperature region above 88°C the contribution of the fast relaxation process decreases with increasing

temperature, and it vanishes in SmC* phase (Fig.1c, curve 1). This fact allows us to assign the fast relaxation process to fluctuations of the antiferroelectric order parameter^[7,10] (or optic-like phason^[15]).

Two low frequency relaxation processes, which appear in all smectic phases, including the SmA, have been assigned to hydrodynamic fluctuation modes due to the existence of one-dimensional modulation of smectic layers (stripe-domains)^[17].

The fast relaxation process exists in the temperature range corresponding to antiferroelectric/ ferroelectric phases in bulk as well as in 2000 Å pores. Therefore one can conclude that the antiferroelectric/ ferroelectric phases appear in 2000 Å pores as well as in bulk.

In 200 Å pores the fast relaxation process appears only in the SmC_A phase. This means that the ferroelectric phases do not appear in 200 Å pores. Combining these facts we can conclude that antiferroelectric order in 200 Å pores does not appear, and the ferroelectric phases are not formed. The reasons for that will be discussed below.

Additionally, the analysis of the temperature dependencies of relaxation times shows^[16] that the temperatures of the SmC_A – SmC_γ phase transition in bulk and in pores are of comparable value, whereas the temperature of the SmC* - SmA phase transition in 2000 Å pores is about 3°C lower and in 200 Å porous sample is about 4°C lower than in the bulk sample.

DISCUSSION

The ferroelectric phases exist in 2000 Å pores and do not exist in 200 Å pores. This can be understood on the basis of the Ising model, which uses a free energy of the form^[13]:

$$F = -\sum_i HS_i + \frac{1}{2} \sum_{i,j} J(i-j) \left(\frac{S_i + 1}{2} \right) \left(\frac{S_j + 1}{2} \right)$$

where the spin $S_i = 1(-1)$ applies to ferroelectric (antiferroelectric) type ordering between two neighboring smectic layers.

According to this model the existence of ferrielectric phases is explained by long range electrostatic interactions between local spontaneous polarizations (spins) of smectic layers, while the antiferroelectric and ferroelectric phases is accounted for by the first term of the equation. The stability of the ferrielectric phases depends on the strength of the potential $J(i-j)$, *i.e.*, how rapidly it varies with the distance between smectic layers. In the bulk liquid crystal the smectic layer thickness ($L \approx 30 \text{ \AA}$) is much smaller than the dimensions of the area of smectic layers, giving an interaction along an infinite number of infinite (in extent) smectic layers. In sufficiently small cylindrical pores (200 \AA), the area of the smectic layer becomes commensurate with the smectic layer thickness, the electric field of the spontaneous polarization of smectic layers decays more rapidly (as function of distance) than in bulk and in 2000 \AA pores limiting the number of interacting smectic layers. In other words, the shrinking of the area of smectic layers shortens the electrostatic interaction between local spontaneous polarizations in the direction of smectic layer normal. In this case (for 200 \AA pores) the electrostatic interactions are not strong enough to form ferrielectric phases.

The observed behavior in cylindrical pores is extremely rich, and provides an important platform for better understanding the nature of ferrielectric phases in antiferroelectric liquid crystalline materials.

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